



Lecture 12: Conservation Laws in the Atmosphere (II)

Prof. Seon K. Park (Ewha Womans Univ.)

Prof. Claudio Cassardo (Univ. of Torino)



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Conservation Laws

- Why conservation laws?
- A system of governing equations includes the conservation of mass (air and water vapor, separately), momentum and energy.
- The **governing equations** of a mathematical model describe how the values of the unknown variables (i.e., the dependent variables) change when one or more of the known (i.e., independent) variables change.

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Conservation Laws

- **Governing equations are based on physical principles derived from the conservation of**
 - Mass → (Mass) Continuity Equ. → ρ
 - Momentum → Equ. of Motion (or Momentum Equ.) → $\mathbf{v}=(u,v,w)$
 - Energy (1st Law of Thermodynamics) → Thermodynamic Equ. (or Temperature Equ.) → T or θ
 - Water Mass → Water Vapor Budget (Moisture) Equ. → q_v, \dots, q_h
 - Ideal Gas Law → Equ. of State (General Gas Equ.) → p

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Conservation of Mass

- **We exclude water vapor, for the moment. ← Phase change**
- **Mass is neither created nor destroyed.**
- **The density comes in here instead of mass (mass per volume):**
 - If the density at a location increases, then the density has to be transferred from some other place to this location.
 - Thus, the total change of density in a moving air parcel is proportional to the divergence of the air stream in the volume.

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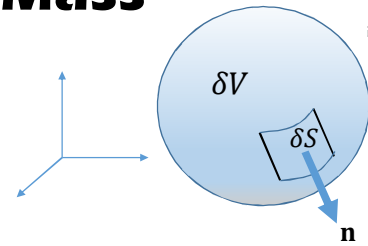
Conservation of Mass

- Lagrangian view:

$$\frac{DM}{Dt} \equiv 0 = \frac{D(\rho V)}{Dt}$$

$$\frac{D\rho}{Dt} + \frac{\rho}{V} \frac{DV}{Dt} = \boxed{\frac{D\rho}{Dt} + \rho \nabla \cdot \mathbf{v} = 0}$$

Continuity Equation
(velocity divergence form)



$$\iiint_V \nabla \cdot \mathbf{Q} dV = \iint_S \mathbf{Q} \cdot \mathbf{n} dS = \iint_S \mathbf{Q} \cdot d\mathbf{s}$$

← Divergence theorem

Practice: Using the divergence theorem, show that

$$\frac{1}{V} \frac{DV}{Dt} = \nabla \cdot \mathbf{v}$$

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Conservation of Mass

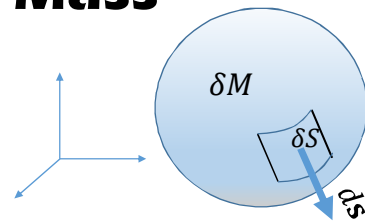
- Lagrangian view

- Considers the mass δM moving together with the air flux at the internal of a material surface δS

$$\frac{DM}{Dt} \equiv 0 = \frac{D(\rho V)}{Dt}$$

$$\frac{D\rho}{Dt} + \frac{\rho}{V} \frac{DV}{Dt} = \boxed{\frac{D\rho}{Dt} + \rho \nabla \cdot \mathbf{v} = 0}$$

Continuity Equation



$$\frac{D\delta M}{Dt} \equiv 0 = \frac{D(\rho \delta V)}{Dt} = \delta V \frac{D\rho}{Dt} + \rho \frac{D\delta V}{Dt} \quad (1)$$

Practice: From (1), derive the continuity equation.

$$\delta V \frac{D\rho}{Dt} + \boxed{\rho \frac{D\delta V}{Dt}} = 0 = \delta V \frac{D\rho}{Dt} + \boxed{\delta V \rho (\nabla \cdot \mathbf{v})} \Rightarrow \frac{D\rho}{Dt} + \rho (\nabla \cdot \mathbf{v}) = 0$$

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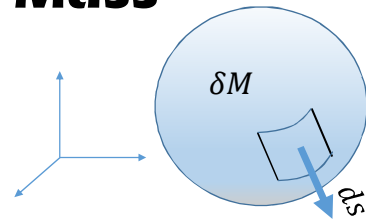
Conservation of Mass

- Lagrangian view:

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Continuity Equation



$$\frac{D\delta M}{Dt} \equiv 0 = \frac{D(\rho \delta V)}{Dt} = \delta V \frac{D\rho}{Dt} + \rho \frac{D\delta V}{Dt}$$

$$\frac{D\delta V}{Dt} = \frac{D}{Dt}(\delta x \delta y \delta z) = \delta y \delta z \frac{D\delta x}{Dt} + \delta x \delta z \frac{D\delta y}{Dt} + \delta x \delta y \frac{D\delta z}{Dt}$$

$$\frac{1}{\delta q} \frac{D\delta q}{Dt} = \frac{\partial}{\partial q} \frac{Dq}{Dt}$$

$$\frac{D\delta V}{Dt} = \delta x \delta y \delta z \left(\left(\frac{\partial}{\partial x} \frac{Dx}{Dt} \right) + \left(\frac{\partial}{\partial y} \frac{Dy}{Dt} \right) + \left(\frac{\partial}{\partial z} \frac{Dz}{Dt} \right) \right) = \delta V \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right)$$

$$\delta V \frac{D\rho}{Dt} + \rho \frac{D\delta V}{Dt} = 0 = \delta V \frac{D\rho}{Dt} + \delta V \rho (\nabla \cdot \mathbf{v}) \Rightarrow \frac{D\rho}{Dt} + \rho (\nabla \cdot \mathbf{v}) = 0$$

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Conservation of Mass

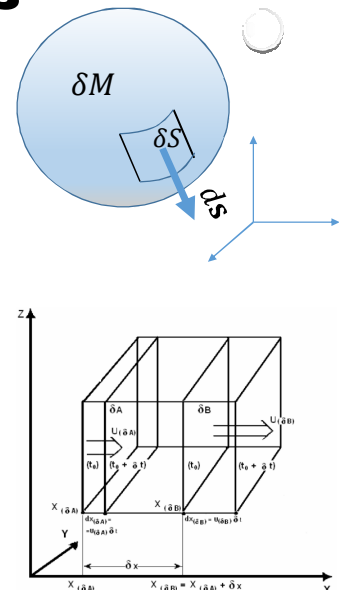
- Lagrangian view: Summary

- Consider a mass δM , of volume $\delta V = \delta x \delta y \delta z$, moving (with δM constant) with the flow
- All particles in δV enclosed by δS never abandon δV remaining forever in δS by definition
- Thus, as $\delta M = \rho \delta V$ is conserved during the motion (despite both ρ and δV can individually vary):

$$\begin{aligned} \frac{1}{\delta M} \frac{D(\delta M)}{Dt} &= \frac{1}{\rho \delta V} \frac{D(\rho \delta V)}{Dt} = \frac{1}{\rho \delta V} \left[\rho \frac{D(\delta V)}{Dt} + \delta V \frac{D\rho}{Dt} \right] = \frac{1}{\rho} \frac{D\rho}{Dt} + \frac{1}{\delta V} \frac{D(\delta V)}{Dt} \\ &= \frac{1}{\rho} \frac{D\rho}{Dt} + \frac{1}{\delta x} \frac{D(\delta x)}{Dt} + \frac{1}{\delta y} \frac{D(\delta y)}{Dt} + \frac{1}{\delta z} \frac{D(\delta z)}{Dt} \end{aligned}$$

- Being $\delta u = D(\delta x)/Dt$ and so on, and taking the limit for $\delta V \rightarrow 0$:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} + \frac{1}{\rho} \frac{D\rho}{Dt} = \nabla \cdot \mathbf{v} + \frac{1}{\rho} \frac{D\rho}{Dt} = 0$$

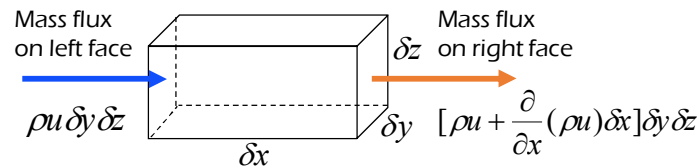


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Conservation of Mass

- **Eulerian view:**

- Considers the entering and leaving of mass fluxes through the surfaces of a cube whose walls are permeable to the air passage



Net increase of mass per unit volume in x -direction: $-\frac{\partial}{\partial x}(\rho u)$

Net increase of mass per unit volume in y -direction: $-\frac{\partial}{\partial y}(\rho v)$

Net increase of mass per unit volume in z -direction: $-\frac{\partial}{\partial z}(\rho w)$

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Conservation of Mass

- **Eulerian view:**

Thus, the net increase of mass per unit volume

$$= -\left(\frac{\partial}{\partial x}(\rho u) + \frac{\partial}{\partial y}(\rho v) + \frac{\partial}{\partial z}(\rho w)\right) = -\nabla \cdot (\rho \mathbf{v})$$

$$= \text{rate of change of density} = \frac{\partial \rho}{\partial t}$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

Continuity Equation
(mass divergence form)

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Conservation of Mass

- **Eulerian view: Summary**

- Mass flux: $\mathbf{j} = \rho \mathbf{v}$
- Mass balance of a fluid volume dV with mass $dM = \rho dV$:

$$\frac{\partial}{\partial t} \int_V \rho dV = - \oint_S \mathbf{j} \cdot d\mathbf{s} = - \oint_S \rho \mathbf{v} \cdot d\mathbf{s}$$

- Exchanging time derivative and integral on the left side and applying the divergence theorem to the right side yields:

- From which: $\int_V \frac{\partial \rho}{\partial t} dV = - \int_V \nabla \cdot (\rho \mathbf{v}) dV \Rightarrow \int_V \left[\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) \right] dV = 0$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

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Conservation of Mass

- **(Mass) Continuity Equation:**

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = \overbrace{\frac{\partial \rho}{\partial t} + \mathbf{v} \cdot \nabla \rho}^{\equiv \frac{D\rho}{Dt}} + \rho \nabla \cdot \mathbf{v} = 0$$

$$\Rightarrow \frac{D\rho}{Dt} + \rho \nabla \cdot \mathbf{v} = 0$$

$$\Rightarrow \frac{1}{\rho} \frac{D\rho}{Dt} + \nabla \cdot \mathbf{v} = 0$$

$$\Rightarrow \frac{D}{Dt} (\ln \rho) + \nabla \cdot \mathbf{v} = 0$$

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Conservation of Mass

- **(Mass) Continuity Equation:**

- Horizontal pressure and density gradients are usually very small (but not in the vertical, see hydrostatic equation)
- In case of an incompressible fluid, we have $\partial\rho/\partial t = 0$ and $\nabla\rho = 0$; thus, from the continuity equation: $\nabla \cdot \mathbf{v} = 0$
 - That is, the velocity field is divergence-free.
- In this case, a horizontal divergence or convergence ($\nabla_H = \partial u/\partial x + \partial v/\partial y$) is usually coincident with vertical air movement ($\partial w/\partial z$).

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Divergence vs. Advection

- **Divergence:** $\nabla \cdot \mathbf{v}$

$$\begin{aligned}\nabla \cdot \mathbf{v} &= \left(\frac{\partial}{\partial x} \hat{\mathbf{i}} + \frac{\partial}{\partial y} \hat{\mathbf{j}} + \frac{\partial}{\partial z} \hat{\mathbf{k}} \right) \cdot (u\hat{\mathbf{i}} + v\hat{\mathbf{j}} + w\hat{\mathbf{k}}) \\ &= \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \quad \Rightarrow \quad \nabla \cdot (\rho \mathbf{v}) = \frac{\partial(\rho u)}{\partial x} + \frac{\partial(\rho v)}{\partial y} + \frac{\partial(\rho w)}{\partial z}\end{aligned}$$

- **Advection:** $\mathbf{v} \cdot \nabla$

$$\begin{aligned}\mathbf{v} \cdot \nabla &= (u\hat{\mathbf{i}} + v\hat{\mathbf{j}} + w\hat{\mathbf{k}}) \cdot \left(\frac{\partial}{\partial x} \hat{\mathbf{i}} + \frac{\partial}{\partial y} \hat{\mathbf{j}} + \frac{\partial}{\partial z} \hat{\mathbf{k}} \right) \\ &= \left(u \frac{\partial}{\partial x} + v \frac{\partial}{\partial y} + w \frac{\partial}{\partial z} \right) \quad \Rightarrow \quad \mathbf{v} \cdot \nabla \rho = u \frac{\partial \rho}{\partial x} + v \frac{\partial \rho}{\partial y} + w \frac{\partial \rho}{\partial z}\end{aligned}$$

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Conservation of Momentum

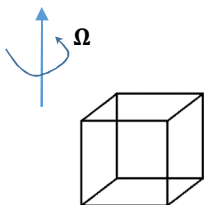
- **Newton's Laws of Motion**

- The momentum of a body remains constant unless the body is acted upon by a net force → **Conservation of Momentum**
- The rate of change of momentum of a body is proportional to the net force acting on the body and is in the same direction of the net force
→ $\mathbf{F} = m\mathbf{a}$
- For every net force acting on a body, there is a corresponding force of the same magnitude exerted by the body in the opposite direction.

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Conservation of Momentum

- **Newton's Laws of Motion**



$$\mathbf{F} = m\mathbf{a} \Rightarrow \frac{\mathbf{F}}{m} = \mathbf{a}$$

$$\mathbf{a}_{\text{net}} = \frac{\mathbf{F}_{\text{net}}}{m} = \frac{\sum_i \mathbf{F}_i}{m} =$$

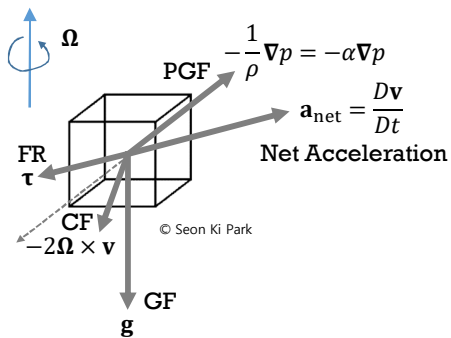
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Conservation of Momentum

- Newton's Laws of Motion

$$\mathbf{F} = m\mathbf{a} \Rightarrow \frac{\mathbf{F}}{m} = \mathbf{a}$$

$$\mathbf{a}_{\text{net}} = \frac{\mathbf{F}_{\text{net}}}{m} = \frac{\sum_i \mathbf{F}_i}{m} = \mathbf{a}_{\text{PGF}} + \mathbf{a}_{\text{CF}} + \mathbf{a}_{\text{GF}} + \mathbf{a}_{\text{FR}}$$

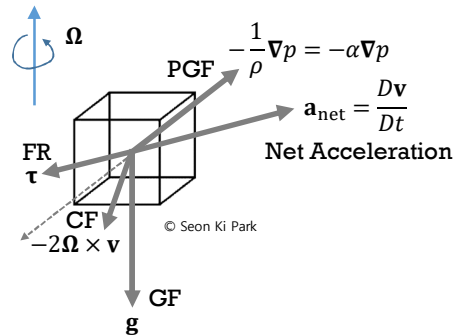


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Conservation of Momentum

- Newton's Laws of Motion

$$\mathbf{F} = m\mathbf{a} \Rightarrow \frac{\mathbf{F}}{m} = \mathbf{a}$$



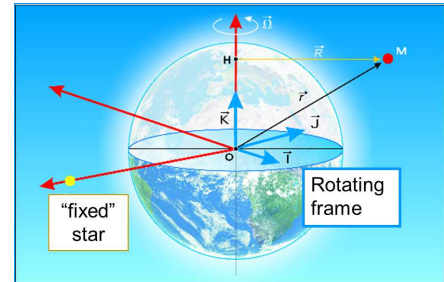
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Conservation of Momentum

- Newton's Laws of Motion → Equ. of Motion (Momentum Equ.)

$$\frac{D\mathbf{v}}{Dt} = \underbrace{-2\boldsymbol{\Omega} \times \mathbf{v}}_{\text{Coriolis acceleration}} + \underbrace{-\boldsymbol{\Omega} \times (\boldsymbol{\Omega} \times \mathbf{r})}_{\text{centrifugal acceleration}} + \underbrace{\mathbf{F}/m}_{\text{acceleration by "absolute" forces}}$$

$$\frac{\mathbf{F}}{m} = \underbrace{(\alpha = 1/\rho) \underbrace{-\alpha \nabla p}_{\text{acceleration by press. grad. force}}}_{\text{acceleration by press. grad. force}} + \underbrace{\mathbf{g}^*}_{\text{acceleration by gravitational force}} + \underbrace{\boldsymbol{\tau}}_{\text{acceleration by frictional force}}$$



(Malardel, 2010; ECMWF Lecture)

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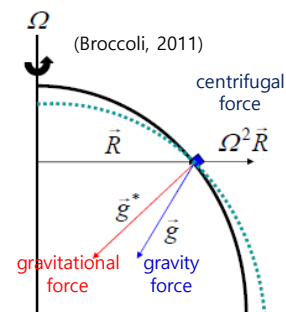
Conservation of Momentum

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$$\frac{\mathbf{F}}{m} = \underbrace{(\alpha = 1/\rho) \underbrace{-\alpha \nabla p}_{\text{acceleration by press. grad. force}}}_{\text{acceleration by press. grad. force}} + \underbrace{\mathbf{g}^*}_{\text{acceleration by gravitational force}} + \underbrace{\boldsymbol{\tau}}_{\text{acceleration by frictional force}}$$

$$\frac{D\mathbf{v}}{Dt} = -\alpha \nabla p - 2\boldsymbol{\Omega} \times \mathbf{v} + \underbrace{-\boldsymbol{\Omega} \times (\boldsymbol{\Omega} \times \mathbf{r}) + \mathbf{g}^*}_{-\mathbf{g}\hat{\mathbf{k}}: \text{gravity accel.}} + \boldsymbol{\tau}$$

$$\frac{D\mathbf{v}}{Dt} = -\alpha \nabla p - 2\boldsymbol{\Omega} \times \mathbf{v} - g\hat{\mathbf{k}} + \boldsymbol{\tau} \quad \leftarrow \text{Equation of Motion or Momentum Equation}$$



The earth has adjusted to this balance of forces by taking the shape of a spheroid with an equatorial bulge.

(Malardel, 2010; ECMWF Lecture)

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Conservation of Energy

• 1st Law of Thermodynamics

- Energy is neither created nor destroyed. → **Conservation of Energy**
- Change in internal energy (U) is due to a combination of heat (Q) added to the system and work done by the system (W).

$$\Delta U = \Delta Q - \Delta W$$

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Conservation of Energy

• 1st Law of Thermodynamics

rate of work done by gas on its surroundings by compression/expansion

$$\frac{dU}{dt} = \dot{Q} - \dot{W}$$

rate of heat exchange with the surroundings

For a perfect gas, $U = c_v T$ (c_v = specific heat at constant volume),

$$\dot{W} = p \frac{d\alpha}{dt} \quad (\alpha = 1/\rho)$$

$$c_v \frac{dT}{dt} + p \frac{d\alpha}{dt} = \dot{Q}$$

diabatic heating rate

rate of change of IE

conversion between thermal and mechanical E

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Conservation of Energy

• Thermodynamic Equation:

$$c_v \frac{dT}{dt} + p \frac{d\alpha}{dt} = \dot{Q}$$

$$c_p \frac{dT}{dt} - \alpha \frac{dp}{dt} = \dot{Q}$$

$\dot{Q} = 0$ for adiabatic process

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Conservation of Energy

• Thermodynamic Equation:

$$c_p \frac{dT}{dt} - \alpha \frac{dp}{dt} = \dot{Q}$$

$$c_p \frac{dT}{dt} = \dot{Q} + \alpha \left(\frac{\partial p}{\partial t} + \mathbf{v}_H \cdot \nabla p + w \frac{\partial p}{\partial z} \right) = \dot{Q} + \alpha \left(\overset{\text{small}}{\cancel{\frac{\partial p}{\partial t}}} + \mathbf{v}_H \cdot \overset{\text{small}}{\cancel{\nabla p}} \right) - gw$$

$$\frac{dT}{dt} = \frac{\dot{Q}}{c_p} - \frac{g}{c_p} w = \frac{\partial T}{\partial t} + \mathbf{v}_H \cdot \nabla T + w \frac{\partial T}{\partial z}$$

$$\frac{\partial T}{\partial t} = \frac{\dot{Q}}{c_p} - \left(\frac{g}{c_p} + \frac{\partial T}{\partial z} \right) w - \mathbf{v}_H \cdot \nabla T = \frac{\dot{Q}}{c_p} - (\Gamma_d - \Gamma) w - \mathbf{v}_H \cdot \nabla T$$

$\Gamma_d = \frac{g}{c_p}; \Gamma = -\frac{\partial T}{\partial z}$

(Broccoli, 2011)

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Conservation of Energy

- **Thermodynamic Equation:**

$$c_p \frac{dT}{dt} - \alpha \frac{dp}{dt} = \dot{Q}$$

$$\frac{\partial T}{\partial t} = \underbrace{\dot{Q}/c_p}_A - \underbrace{(\Gamma_d - \Gamma)w}_B - \underbrace{\mathbf{v}_H \cdot \nabla T}_C$$

- **A: Diabatic Heating**

1. Sensible heating
2. Latent heating (phase change, i.e., evaporation, condensation)
3. Radiative heating

- **B: Adiabatic Effects**

1. $(\Gamma_d - \Gamma)$ is a measure of stability.
2. Upward motion in a stable atmosphere is a cooling process.

- **C: Horizontal Advection Term**

1. Often this is the largest term

(Broccoli, 2011)

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